

and Ermalov⁷ who concluded from studies of the phosphorescence spectra of mixtures of benzophenone and other donors) and naphthalene in ether plus ethanol at 77°K that there was a transfer of energy between the triplet states of these molecules in the rigid glass solutions.

Our results also suggest a possible explanation of the observations of Ferguson and Tinson⁸ who found that the phosphorescence spectrum of benzophenone in petroleum ether at -180°C is replaced by another at longer wavelengths when the concentration is increased from 10⁻⁴ to 10⁻²M. These workers attributed their results to energy transfer to the triplet state of benzophenone. As we have shown that the triplet state of this molecule reacts very rapidly at 77°K with the components of E.P.A. to form a ketyl radical, it is possible that the second phosphorescence spectrum observed by Ferguson and Tinson comes from the ketyl radical or some resulting compound.

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Sixth Virial Coefficients for Gases of Parallel Hard Lines, Hard Squares, and Hard Cubes

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THE virial equation of state can be written as

$$\beta p = \sum_1^{\infty} B_n \rho^n,$$

where $\beta = 1/kT$, ρ is the number density, and the set of coefficients $\{B_n\}$ are the virial coefficients. The topological graphs are known which permit the calculation of $\{B_n\}$ up to B_7 ,¹ if the related integrals of the f_{ij} products can be accomplished. Because the integrals are formidable for realistic choices of the potential, simpler models have been introduced to effect the

integrals. The simplest nontrivial case is that of parallel hard lines, squares, and cubes.² The hard-cube potential has the property that any integral I_n of a product πf_{ij} can be written as $I_1 |I_1|^{n-1}$, where n is the dimensionality of the interacting figures; this feature of the problem was clearly noted by Zwanzig.³ To compute the set $\{B_n\}$ one need only evaluate the pertinent set of integrals in one dimension; the extension to higher dimensions involves only mechanical effort. To calculate B_2 and B_3 , only one integral must be done. For B_4 , B_5 , and B_6 one needs to accomplish three, 10, and 56 integrals, respectively.

We have succeeded in computing B_6 by a straightforward but tedious method, the details of which will be published. We present in Table I the results for B_6

TABLE I. Virial coefficients for lines, squares, and cubes.

	Lines	Squares	Cubes
B_1	1	1	1
B_2	1	2	4
B_3	1	3	9
B_4	1	11/3	34/3
B_5	1	67/18	455/144
B_6	1	121/40	-2039/108
B_1	1.0000	1.0000	1.0000
B_2	1.0000	1.0000	1.0000
B_3	1.0000	0.7500	0.5625
B_4	1.0000	0.4583	0.1771
B_5	1.0000	0.2326	0.0123
B_6	1.0000	0.0945	-0.0184

and the lower coefficients (previously known), for lines of length σ , squares of area σ^2 , and cubes of volume σ^3 . The first set of values corresponds to the choice of σ^n for the n -dimensional B_2 ; the second set uses $2^{n-1}\sigma^n$ for the corresponding B_2 (both sets are employed in the literature). It is interesting to note that B_6 is negative in three dimensions, a result anticipated by Temperley on the basis of his extension of the Zwanzig calculations to higher dimensions.⁴ At this point the earlier work deserves some further comment:

(1) Tonks⁵ calculated the equation of state for the one-dimensional case, but the fact that Tonks did not employ the cluster expansion prevents any extension of his results to higher dimensions.

(2) Geilikman² obtained B_2 , B_3 , and B_4 for cubes, although, as he later pointed out, the result for B_4 was incorrect.

(3) Zwanzig,³ using the results of Riddell and Uhlenbeck,⁶ computed the first five coefficients for the case described in this letter. Zwanzig agrees that the inequalities in Helmholtz energy for systems of cubes and spheres used in his paper³ to limit the virial coefficients for spheres do not seem sufficient for this purpose.⁷ If his treatment had been rigorous our result would imply a negative B_6 for hard spheres, in contrast to the suggestion of Alder and Wainwright⁸ that both B_6

and B_7 for hard spheres are positive. Whether or not B_n° and B_n^{\square} have the same sign is still an open question.

(4) As noted earlier Temperley found that B_4 was negative in five dimensions and B_5 negative in four dimensions. It is clear that B_5 will eventually become positive again if a sufficient number of dimensions is employed. His later statement⁹ that Ford found B_4 and B_7 as negative for spherical molecules is without foundation, since the Gaussian model employed by Ford does not correspond to any potential.

(5) Domb¹⁰ presented equations of state for "squares" and "cubes" in which B_4 , B_5 , and B_6 were negative. These were presumably approximations derived from an Ising model, but none of the coefficients beyond the first agree with the exact results for hard squares and hard cubes.

We will present the detailed calculation of B_6 in the near future. We are now actively concerned with the calculation of B_7 and the analog of the radial distribution function for such systems.

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Infrared Spectrum of Diimide*

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BANDS belonging to the infrared spectrum of the diimide (N_2H_2) molecule have been observed. This molecule has been previously identified by Foner and Hudson¹ by mass spectrometry as a product of a radio-frequency discharge though gaseous hydrazine. We have used the method of preparation of Foner and Hudson. Our experimental arrangement permits deposition of the discharge products on a cold infrared-transparent window and simultaneous sampling of the

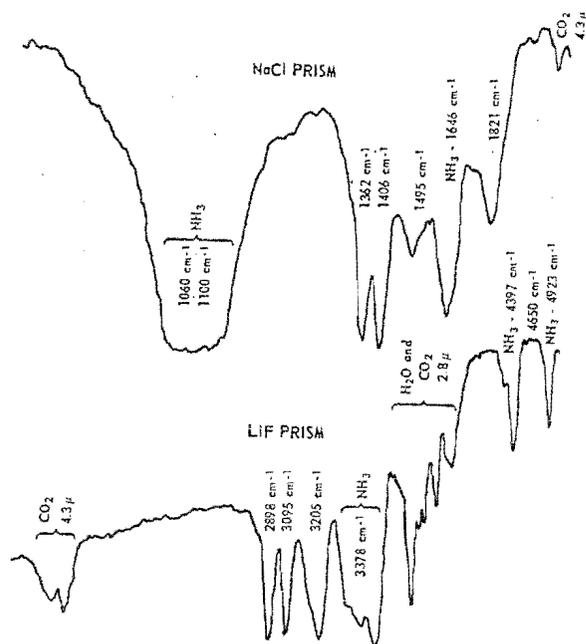


FIG. 1. Infrared spectrum of diimide.

discharge products for analysis on the mass spectrometer. A dry-ice trap immediately follows the discharge to remove hydrazine. The cryostat is very similar to one described by Wagner and Hornig.² It has a clear aperture of 35 mm and fits easily into the sample space of the modified Perkin-Elmer model 12B spectrometer used in these experiments. With liquid nitrogen as coolant, a calcium fluoride window reached a steady temperature of 82°K in 15 min and a sodium chloride window 86°K in a similar time.

If hydrazine vapor is admitted into the evacuated system, it condenses completely in the dry-ice trap and nothing deposits on the cold window of the cryostat. When the discharge is started, the pressure following the dry-ice trap increases to between 100 and 1000 μ , depending on the discharge power. A yellow deposit, which is stable at least 2 to 3 hr at temperatures below 120°K, forms on the cold window.

Analysis of the discharge products with the mass spectrometer shows the presence of nitrogen, ammonia, and diimide. The undecomposed hydrazine is completely collected in the dry-ice trap as indicated by its absence in the mass spectrometer. In some experiments a trace of mass 45 (N_2H_3) was detected. By varying the hydrazine flow rate and the discharge power, it was possible to vary the ratio of diimide to ammonia from below 1% to about 20%.

The spectrum of the deposit as recorded with the NaCl and LiF prisms is shown in Fig. 1. The rock salt region was examined as far as 15 μ , but no bands were found beyond those shown in the tracing. In addition to the bands of solid ammonia,³ the spectrum includes

several known hydrazonium dinitropentoxide bands, attributed to (2898, 3205, and 4650 μ) bands, and a band at 4.3 μ . When the deposit appears.

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* Work of the U.S. Navy (1950), (1951), (1957).